

Fig. 3. Pressure effect on  $T_{\alpha}$  as determined from basic  $C_{33}$  vs. temperature data.  $(C_{33} = \rho V_{[001]}^2)$ 

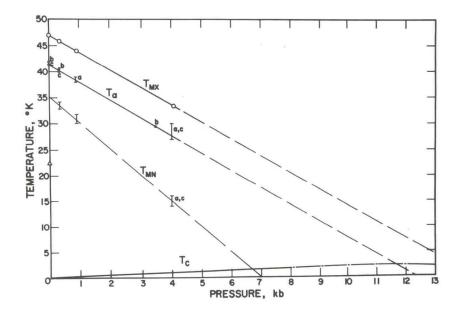


FIG. 4. Variations with hydrostatic pressure of  $T_{\alpha}$  and the temperatures delineating the high acoustic attenuation for  $C_{11}$  mode in uranium. The  $T_c$  vs. pressure curve obtained from Reference 9 to 8 kbar and reference 7 at p>8 kbar.

integral number of specimen waves, V= wave velocity, t= crystal thickness and the subscripted terms represent the basis data at  $77.8^{\circ}$  K.  $T_{\alpha}$  is defined as the temperature at which there occurs

an abrupt change in slope of the  $f/f_{77.8}$  vs. T curves. For the [100] data, Fig. 1,  $T_{\alpha}$  cannot be directly observed because of the very large acoustic attenuation that is associated with the

softness of this mode near the phase transition. The decrease with pressure in  $T_{\alpha}$  and the temperature of high attentuation is, however, clearly indicated, as is the reduction of the hysteresis effects<sup>2</sup> at  $T < T_{\alpha}$ . The data for the [010] direction, Fig. 2, clearly show  $T_{\alpha}$  at two different hydrostatic pressures as well as at 1 bar or 0 psi. The [001] data, Fig. 3, show that unambiguous shift in  $T_{\alpha}$  between 1 bar and 0.35 kbar (5,000 psi). At 4.07 kbar, however, there is a clear change in the character of the data for [001] so that  $T_{\alpha}$  is less well defined.

The data for  $dT_{\alpha}/dP$  are summarized in Fig. 4, where the indicated  $T_{\alpha}$  marked a, b and c refer to [100], [010] and [001], respectively. The points (a) were obtained by interpolation of the [100] data. The decrease in  $T_{\alpha}$  with increasing pressure is very near linear, with  $dT_{\alpha}/dP = -3.4^{\circ} \text{K/kbar}$ . The lines marked  $T_{MX}$  and  $T_{MN}$  delineate the temperature range at a given pressure over which the longitudinal [100] mode signal is lost to attenuation and the pressure range at a given T that this mode velocity decreases with increasing pressure. Finally, the reported superconducting  $T_{\alpha}$  vs. pressure data are reproduced

from reference 9 and 7. We arrive at the conclusion that the maximum  $T_c$  occurs at or near the pressure at which  $T_c = T_\alpha$ . This is a rather different conclusion than is reached from a direct volume to valence relation. It is consistent, however, with the concept that  $T_{c, max}$  occurs when the electronic effects produced by cooling from  $T_{\alpha}$ are completely reversed by ~ 10 kbar pressure or, similarly, ~ 10 kbar pressure is sufficient to retain the bulk superconducting a phase during cooling. It is not, however, necessarily inconsistent with an electron-phonon coupling model. The indications are that  $T_c$ , max occurs where  $\omega_{[100]}$ , the frequency for [100] longitudinal phonons, reaches a minimum value under hydrostatic pressure. In subsequent experiments we will attempt to measure the pressure dependence of the transverse mode velocities at 4°K and thereby estimate the changes in the whole phonon spectrum with pressure at superconducting temperatures.

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